

How to hide a supermassive Black Hole:  
AGN obscuration through dusty infrared dominated flows.

Anton Dorodnitsyn  
NASA GSFC / UMD

# Special Thanks

- Tim Kallman
- G.S. Bisnovaty-Kogan

Dorodnitsyn, Kallman, Bisnovaty-Kogan,  
ApJ, 2012, 747, 8

Dorodnitsyn, Bisnovaty-Kogan, Kallman,  
ApJ, 2011, 741, 29

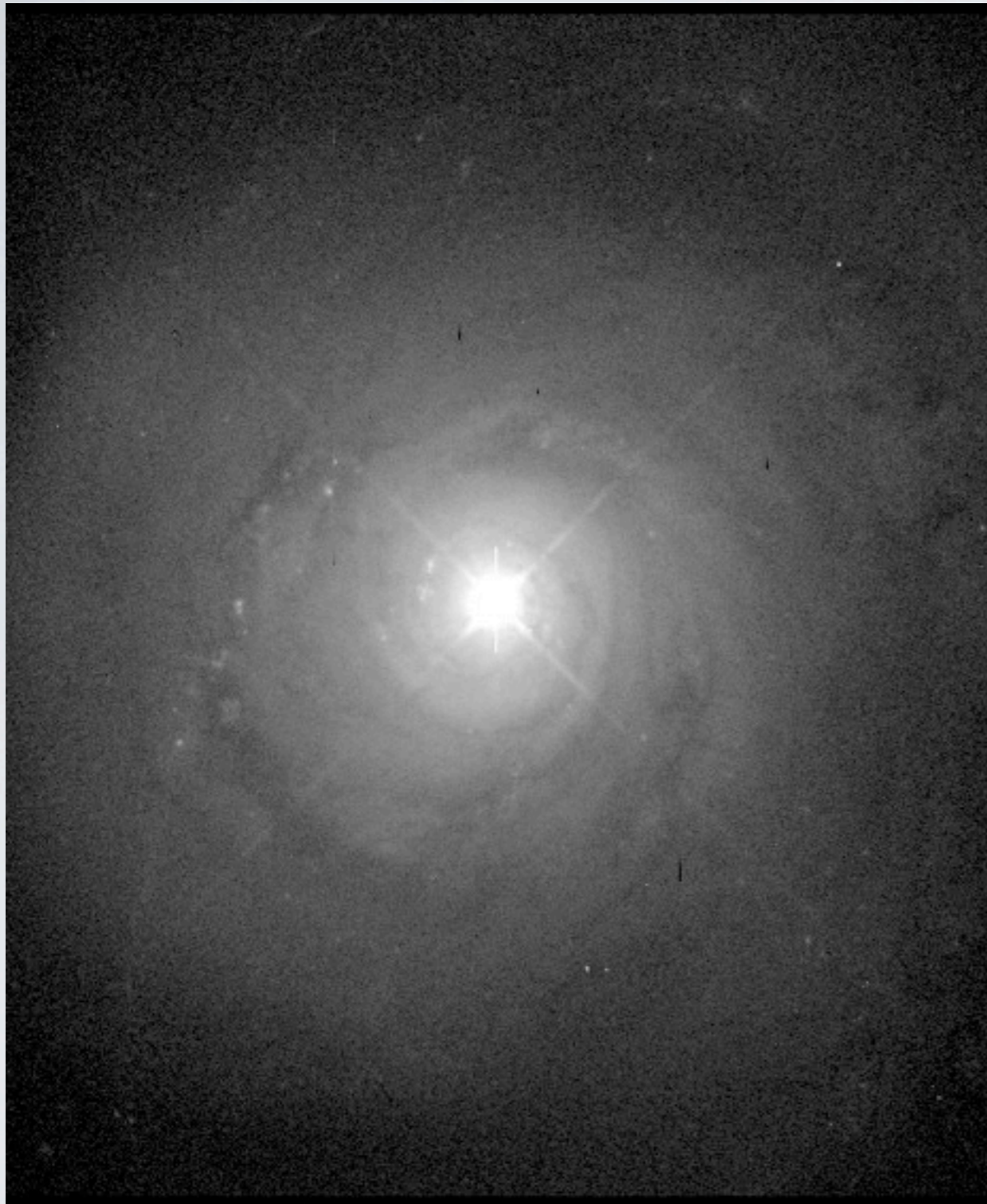
Dorodnitsyn, A., Kallman, T.  
2009, ApJ, 703, 1797

Dorodnitsyn, A., Kallman, T.  
2010, ApJL, 711, 112

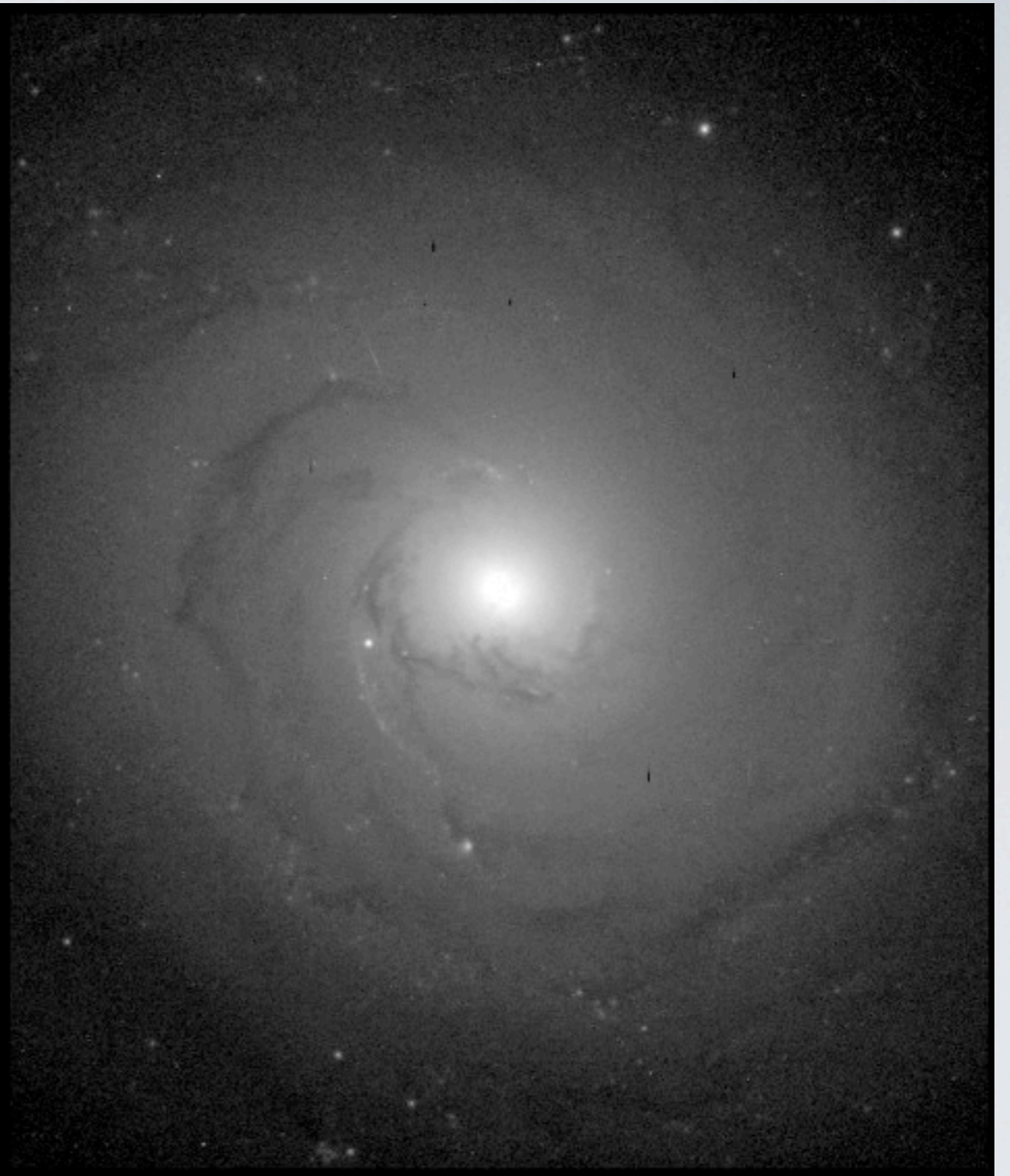
Dorodnitsyn, Kallman, Proga, D. 2008,  
ApJL, 657, 5

Dorodnitsyn, A., Kallman, T., Proga, D.  
2008, ApJ, 687, 97



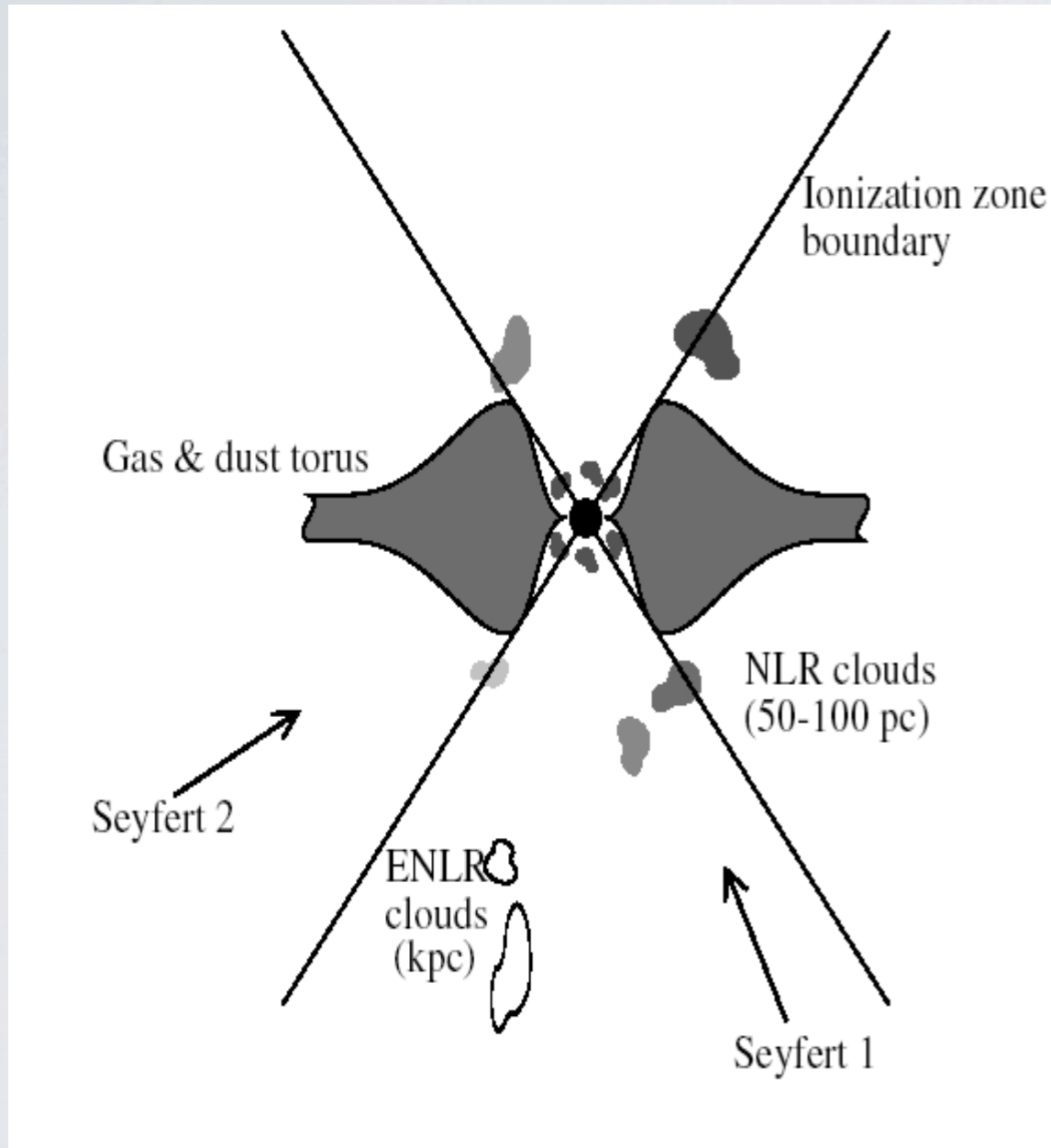
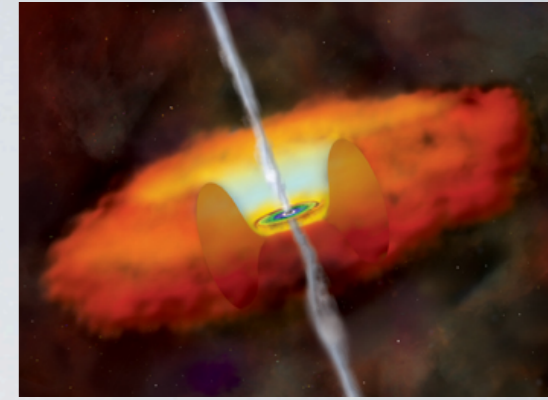


Seyfert 1, galaxy, NGC 5548



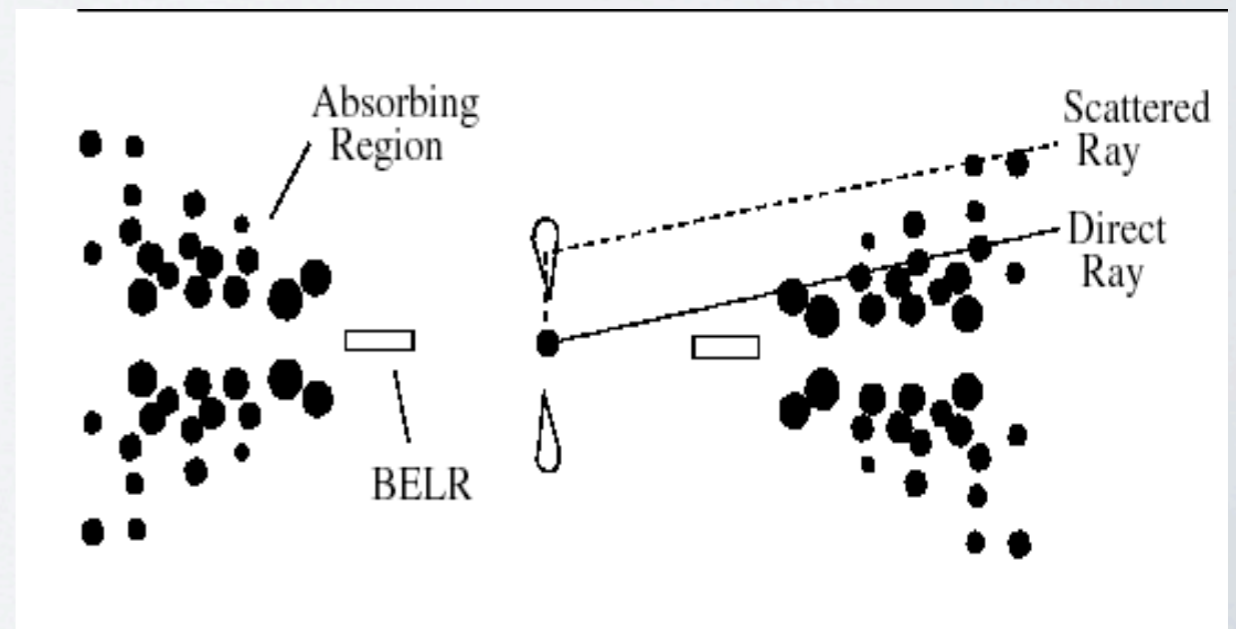
NGC 3277

# AGN UNIFICATION

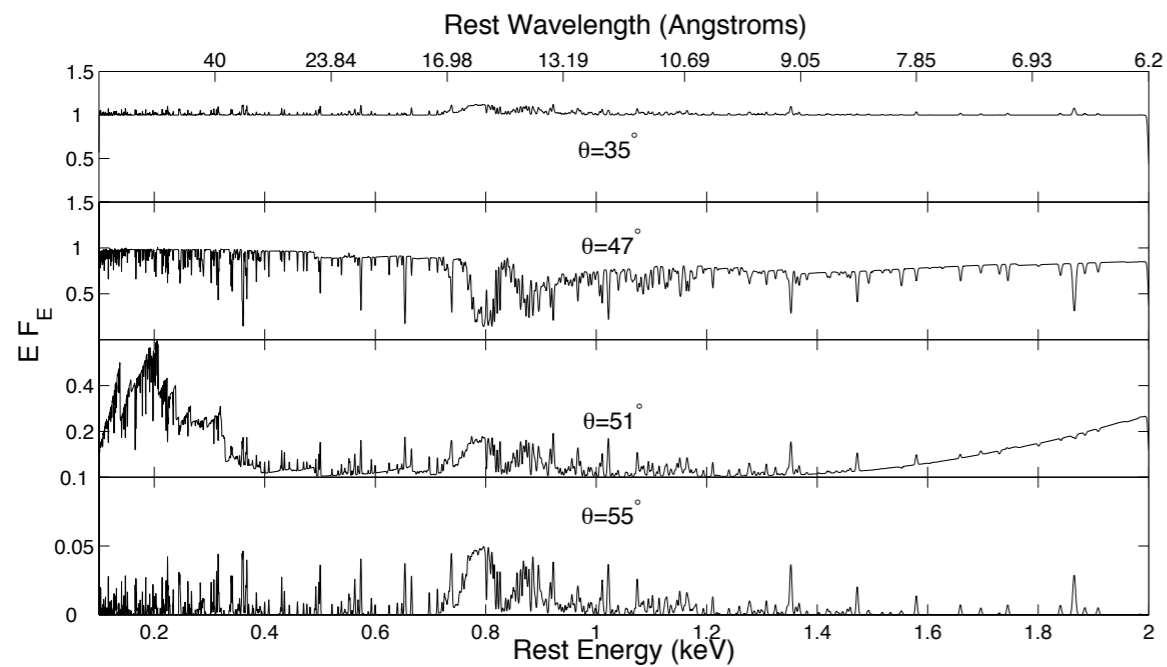


- Seyfert 1: broad emission lines
- Seyfert 2: narrow lines only

So, “type 2” AGN seem to be “type 1” but with no broad lines. Typically, Type 2 are less luminous in optics

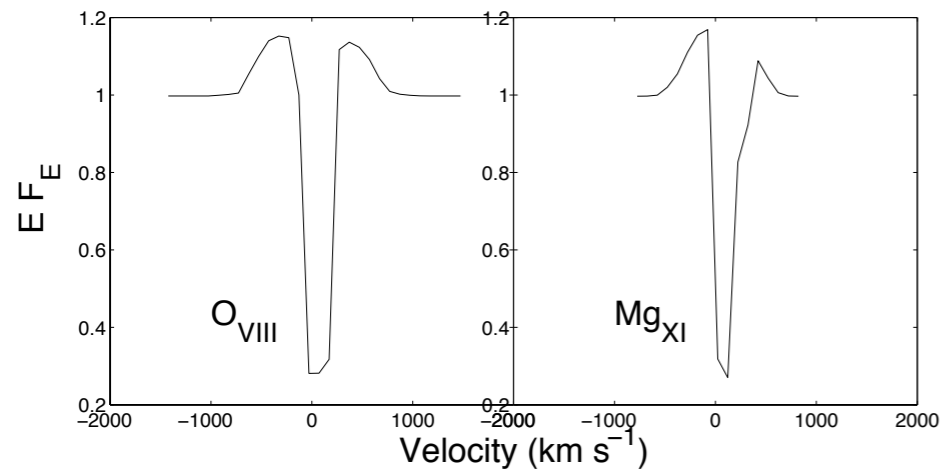






Model X-ray spectra,  
observed as a function of angle.

Dorodnitsyn, A., Kallman, T.  
2009, ApJ, 703, 1797

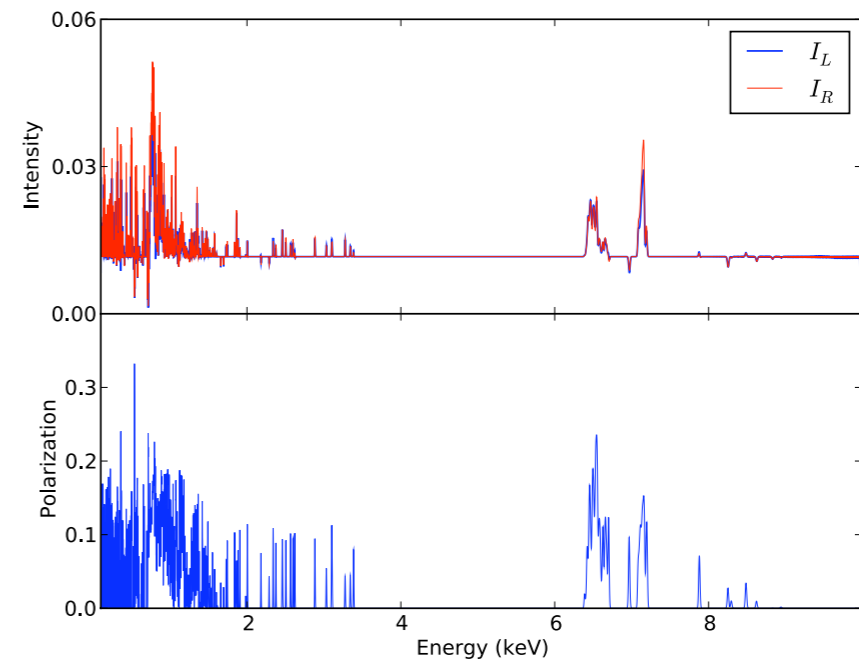


Profiles of individual lines of OVII(left) and MgXI (right).  
Blueshifts correspond to positive velocities, redshifts to negative

## Warm absorbers

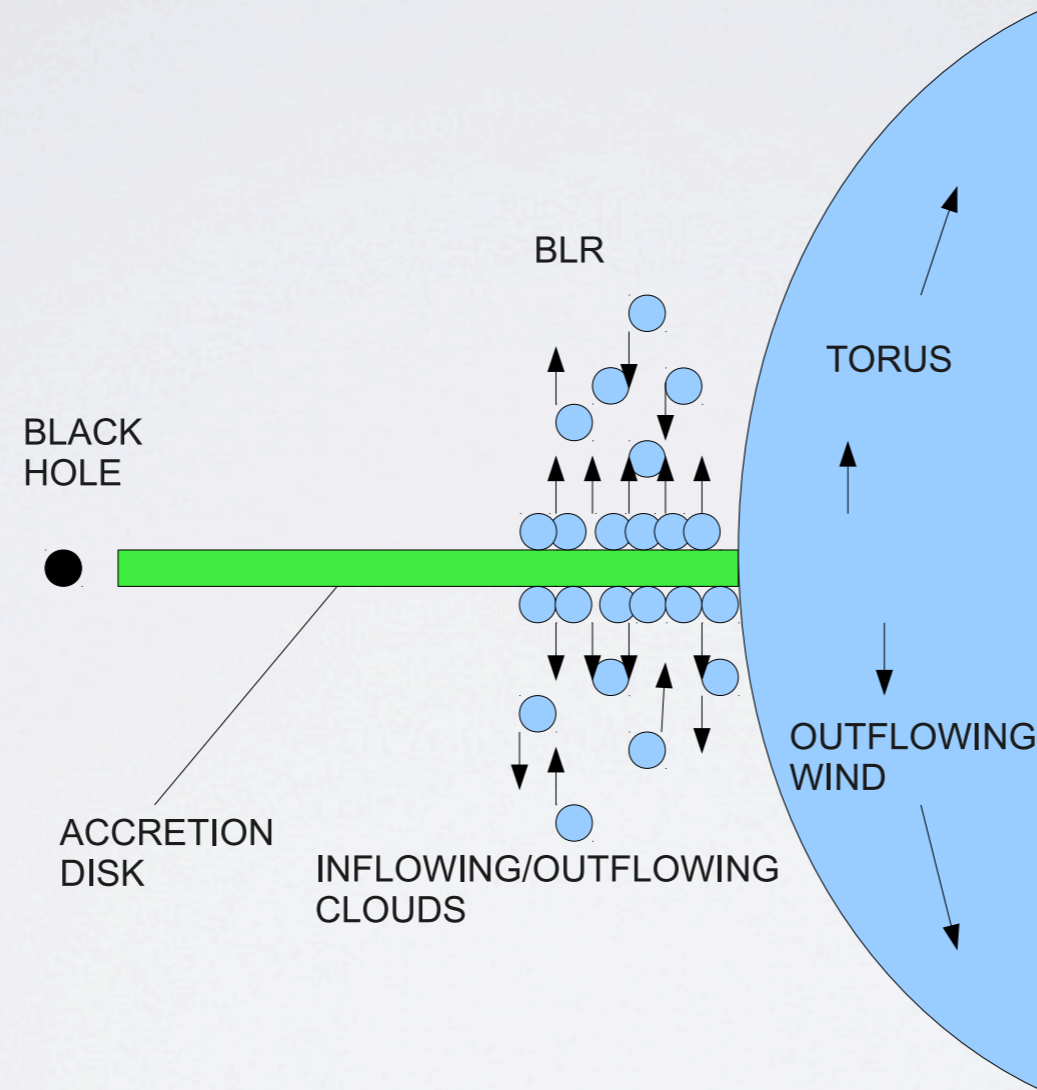
- The grating spectrographs on the X-ray telescopes *Chandra* and *XMM-Newton* provide unprecedented spectral resolution up to  $\sim 10$  keV.
- These show that X-ray spectra obtained from  $\sim$  half of low-red-shift active galactic nuclei (AGN) contain many lines from ions of Fe, Si, S, O, Mg, and Ne, and that these are generally broadened and blueshifted by 100-500 km/s  
Kaspi, S., et al. 2002, ApJ, 574, 643; Steenbrugge, K.C. 2005, A&A, 432, 453
- The presence of X-ray absorbing gas has been confirmed in the majority of AGNs which are bright enough to allow detections.
- There is also a partial correspondence between UV and X-ray absorbers

## X-ray polarimetry of warm absorbers



Dorodnitsyn, A., Kallman, T. 2010, ApJL,  
711,112

# POSSIBLE CONNECTION WITH BLR

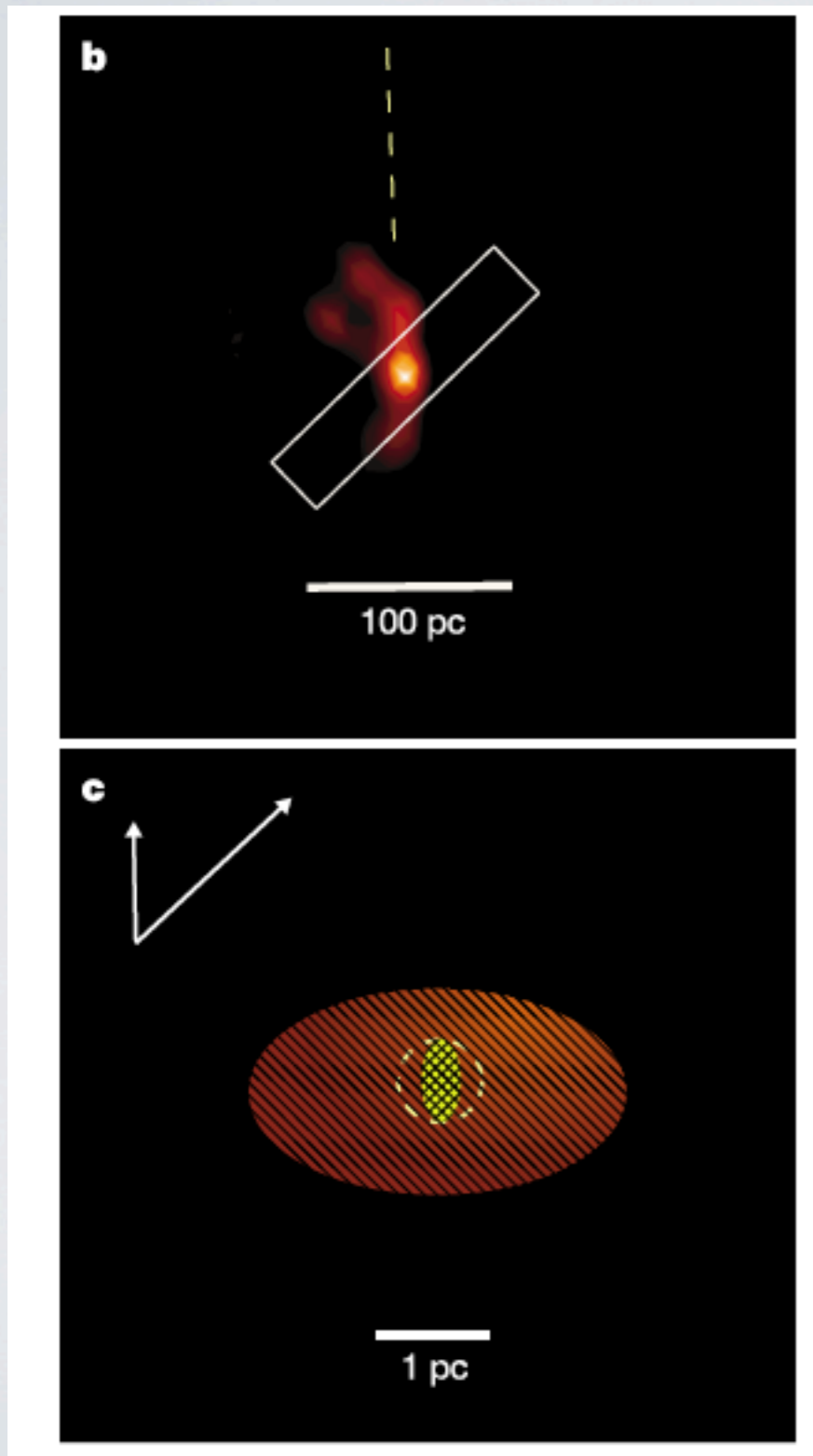




# ALTERNATIVES FOR THE OBSCURATION

- Warps in accretion disk at a pc scale
- Clumps/clouds
- Winds from short distance from the BH
  - Radiation-driven winds
  - MHD (Blandford-Payne -type) flows
- IR-driven

# DIRECT OBSERVATIONS OF THE TORUS



- Orientation effects play a major role in explaining many of the energetic phenomena in AGN
- Dusty clouds form a torus - like structure of the 1 parsec size
- Interferometric mid -infrared observations of NGC 1068 (Seyfert 2) reveals multi-component structure
- Warm (320K) - 2.1 parsec thick and 3.4 parsec in



One of the major problems which should be addressed by a theory of AGN obscuration is how the torus resists collapse into a geometrically thin disk. If the torus is supported by rotation and **gas pressure** then the temperature of such gas should be of the order of the virial temperature:

$$T_{\text{vir,g}} = 2.6 \times 10^5 M_6 / r_{\text{pc}} \text{ K}$$

Clearly, such temperatures cannot be reconciled with the existence of dust.





Comparing the energy density of the X-ray and UV-photons,

$$3.44 \times 10^{-5} M_6 / r_{\text{pc}}^2 \text{ erg cm}^{-3}$$

Comparing the energy density of the X-ray and UV-photons,

$$3.44 \times 10^{-5} M_6 / r_{\text{pc}}^2 \text{ erg cm}^{-3}$$

assuming that the black hole radiates at half of its Eddington luminosity and a 30% covering fraction of the Compton thick portion of the torus



Comparing the energy density of the X-ray and UV-photons,

$$3.44 \times 10^{-5} M_6 / r_{\text{pc}}^2 \text{ erg cm}^{-3}$$

assuming that the black hole radiates at half of its Eddington luminosity and a 30% covering fraction of the Compton thick portion of the torus

with the energy density of infrared radiation, we obtain that a gas-dust temperature of a few x 100K is required if all these X-ray and UV photons are converted to the infrared. A more elaborate treatment shows that if the temperature of the torus is a fraction of

$$T_{\text{vir,r}} \simeq 527 \left( \frac{n/10^7}{r_{\text{pc}}} \right)^{1/4}$$

Comparing the energy density of the X-ray and UV-photons,

$$3.44 \times 10^{-5} M_6 / r_{\text{pc}}^2 \text{ erg cm}^{-3}$$

assuming that the black hole radiates at half of its Eddington luminosity and a 30% covering fraction of the Compton thick portion of the torus

with the energy density of infrared radiation, we obtain that a gas-dust temperature of a few x 100K is required if all these X-ray and UV photons are converted to the infrared. A more elaborate treatment shows that if the temperature of the torus is a fraction of

$$T_{\text{vir,r}} \simeq 527 \left( \frac{n/10^7}{r_{\text{pc}}} \right)^{1/4} \quad \text{the torus thickness will be maintained by IR radiation pressure}$$



Comparing the energy density of the X-ray and UV-photons,

$$3.44 \times 10^{-5} M_6 / r_{\text{pc}}^2 \text{ erg cm}^{-3}$$

assuming that the black hole radiates at half of its Eddington luminosity and a 30% covering fraction of the Compton thick portion of the torus

with the energy density of infrared radiation, we obtain that a gas-dust temperature of a few x 100K is required if all these X-ray and UV photons are converted to the infrared. A more elaborate treatment shows that if the temperature of the torus is a fraction of

$$T_{\text{vir,r}} \simeq 527 \left( \frac{n/10^7}{r_{\text{pc}}} \right)^{1/4} \quad \text{the torus thickness will be maintained by IR radiation pressure}$$

$$\beta = P_g/P \simeq \left( 10^3 \frac{T_3^3}{n_7} + 1 \right)^{-1}$$





$$L_{\text{edd}} = \frac{4\pi cGM_{\text{BH}}}{\kappa_e} = 1.26 \times 10^{44} M_6$$

$$L_{\text{edd}} = \frac{4\pi cGM_{\text{BH}}}{\kappa_e} = 1.26 \times 10^{44} M_6$$

$$L_{\text{c,dust}}^{\text{UV}} = 5 \times 10^{-4} - 0.01 L_{\text{edd}}$$



$$L_{\text{edd}} = \frac{4\pi cGM_{\text{BH}}}{\kappa_e} = 1.26 \times 10^{44} M_6$$

$$L_{\text{c,dust}}^{\text{UV}} = 5 \times 10^{-4} - 0.01 L_{\text{edd}}$$

$$L_c = \frac{4\pi cGM}{\kappa_d} \simeq (0.03 - 0.1) L_{\text{edd}}$$

$$L_{\text{edd}} = \frac{4\pi cGM_{\text{BH}}}{\kappa_e} = 1.26 \times 10^{44} M_6$$

$$L_{\text{c,dust}}^{\text{UV}} = 5 \times 10^{-4} - 0.01 L_{\text{edd}}$$

$$L_c = \frac{4\pi cGM}{\kappa_d} \simeq (0.03 - 0.1) L_{\text{edd}}$$

Purely local,  
no transfer effects!

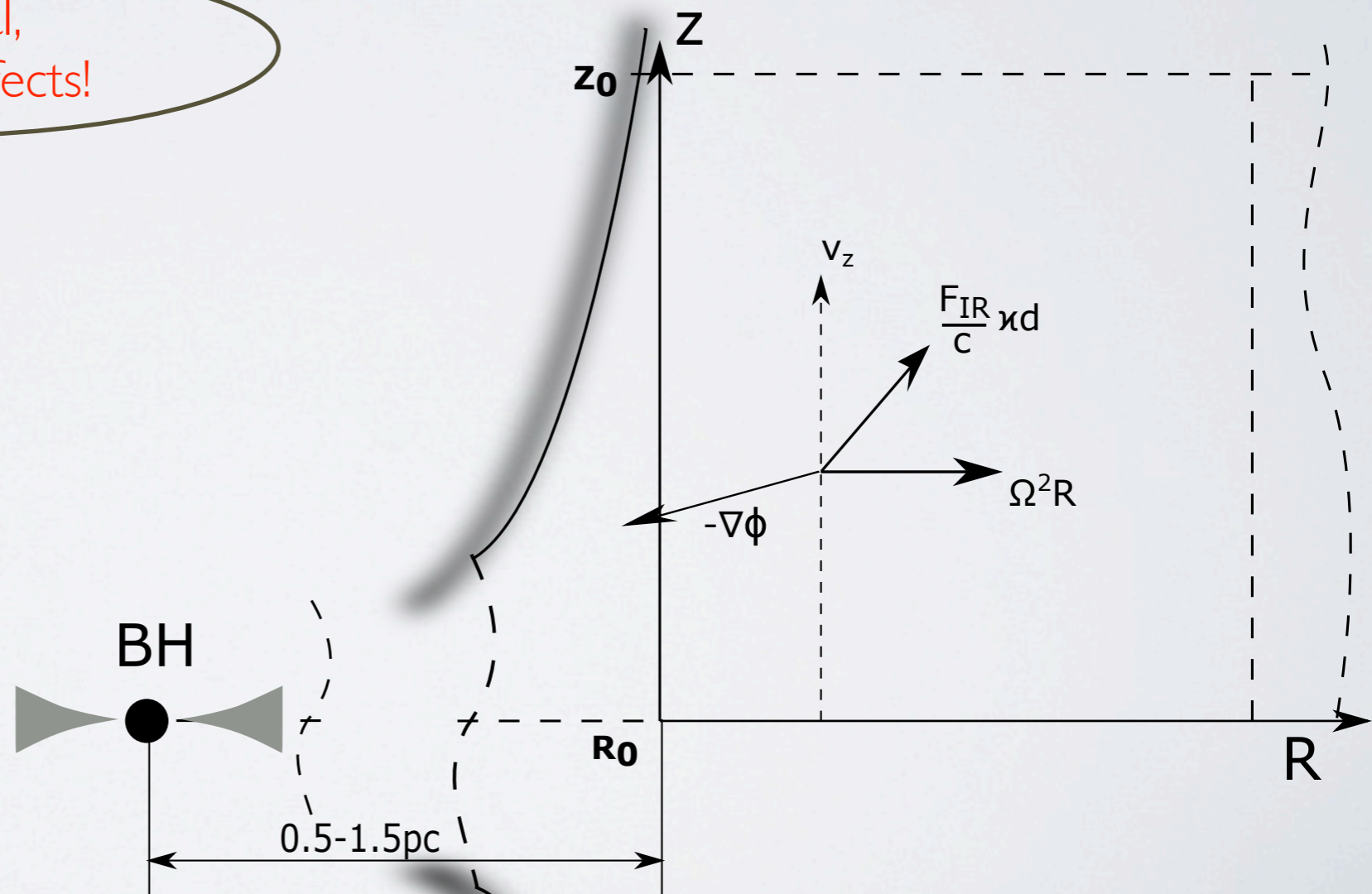


$$L_{\text{edd}} = \frac{4\pi cGM_{\text{BH}}}{\kappa_e} = 1.26 \times 10^{44} M_6$$

$$L_{\text{c,dust}}^{\text{UV}} = 5 \times 10^{-4} - 0.01 L_{\text{edd}}$$

$$L_c = \frac{4\pi cGM}{\kappa_d} \simeq (0.03 - 0.1) L_{\text{edd}}$$

Purely local,  
no transfer effects!



$$L_{\text{eddd}} = \frac{4\pi c G M_{\text{BH}}}{\kappa_e} = 1.26 \times 10^{44} M_6$$

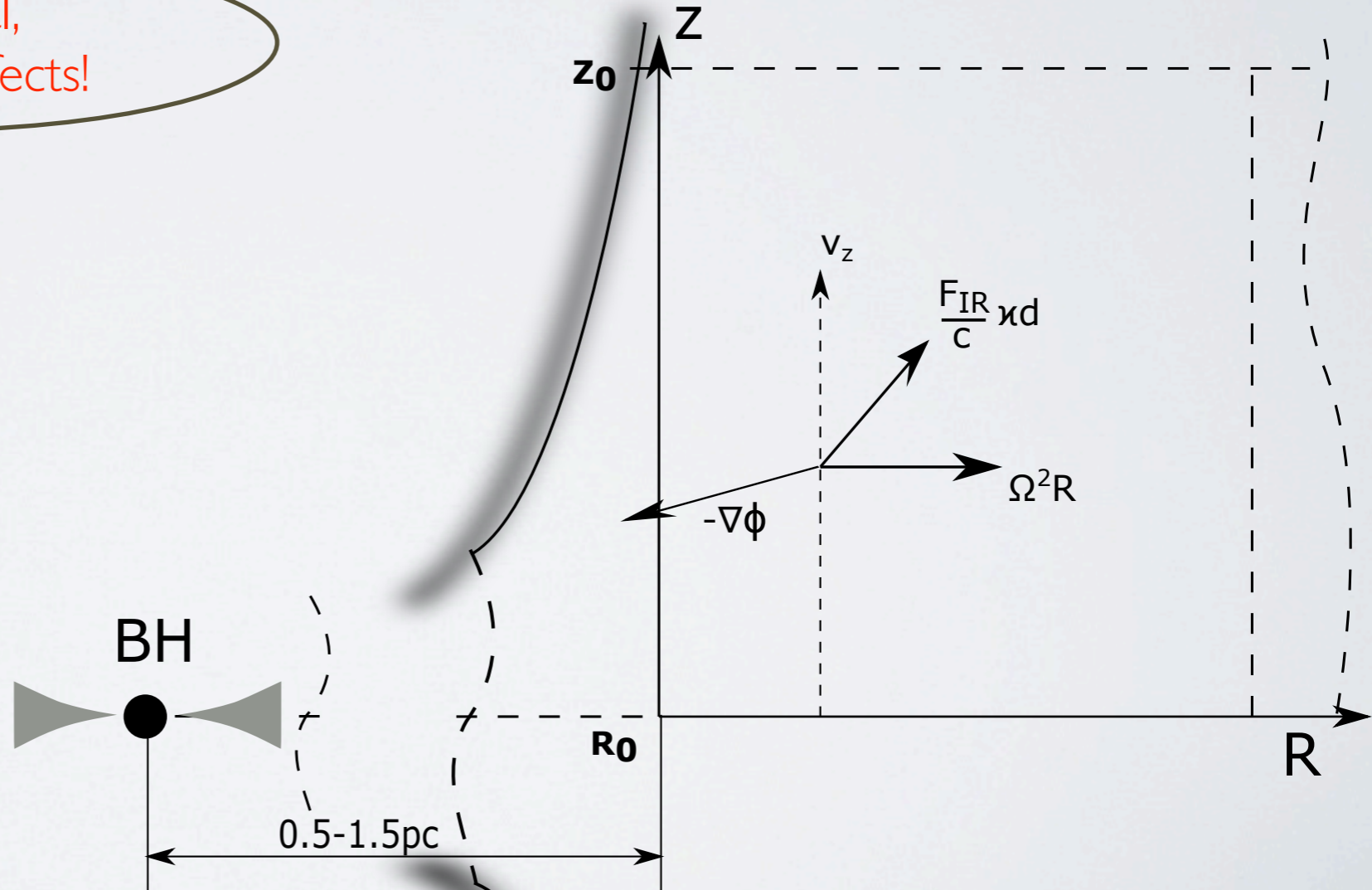
$$L_{\text{c,dust}}^{\text{UV}} = 5 \times 10^{-4} - 0.01 L_{\text{eddd}}$$

$$L_c = \frac{4\pi c G M}{\kappa_d} \simeq (0.03 - 0.1) L_{\text{eddd}}$$

Purely local,  
no transfer effects!

UV is absorbed in a layer

$$\delta l / R_{1\text{pc}} = 1.3 \times 10^{-3} n_7^{-1}$$





$$L_{\text{edd}} = \frac{4\pi c G M_{\text{BH}}}{\kappa_e} = 1.26 \times 10^{44} M_6$$

$$L_{\text{c,dust}}^{\text{UV}} = 5 \times 10^{-4} - 0.01 L_{\text{edd}}$$

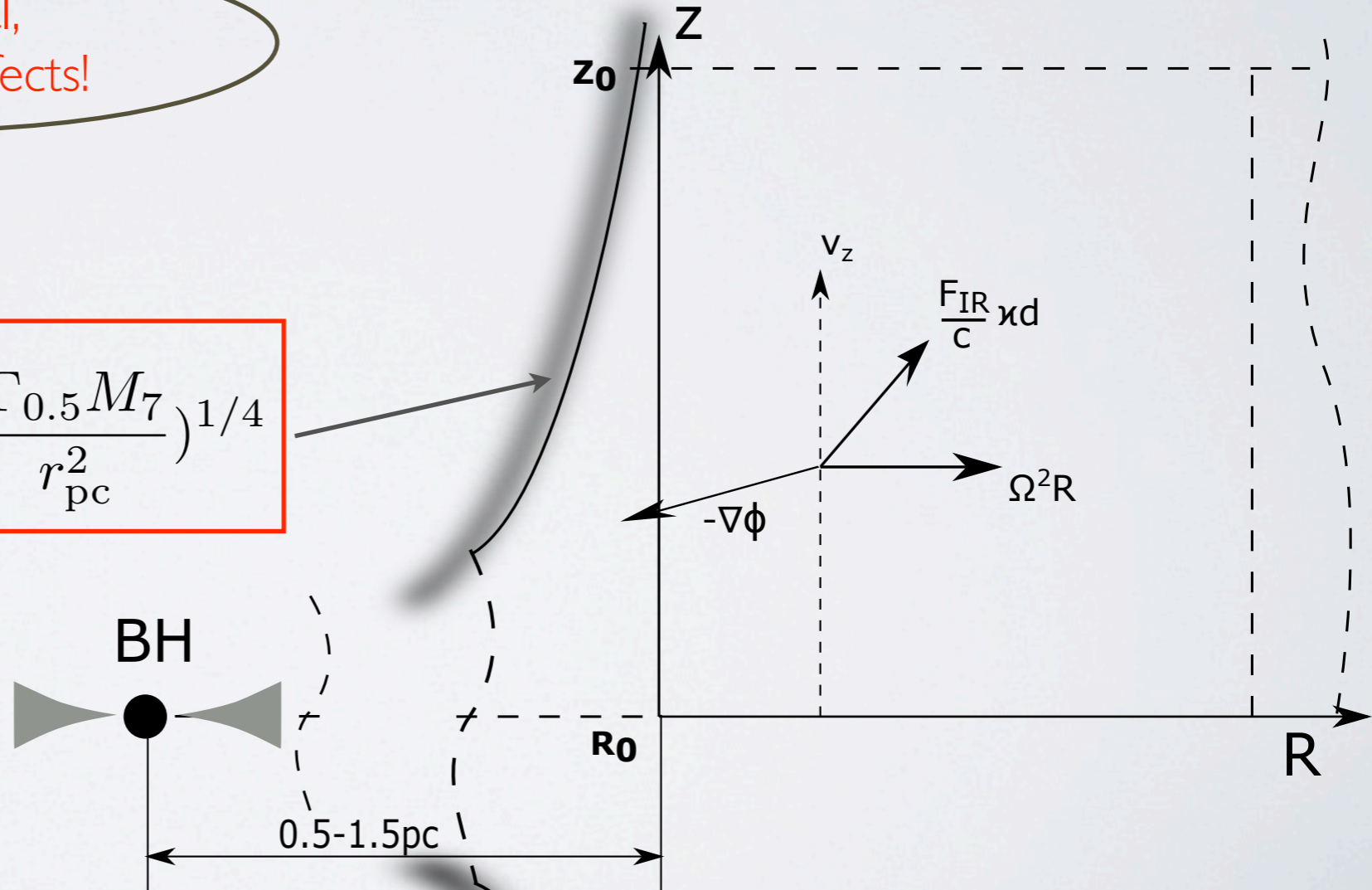
$$L_c = \frac{4\pi c G M}{\kappa_d} \simeq (0.03 - 0.1) L_{\text{edd}}$$

Purely local,  
no transfer effects!

$$T_{\text{eff}} = \left( 4\alpha\Gamma \frac{GM}{\kappa_T a r^2} \right)^{1/4} \simeq 463 \left( \frac{\Gamma_{0.5} M_7}{r_{\text{pc}}^2} \right)^{1/4}$$

UV is absorbed in a layer

$$\delta l / R_{1\text{pc}} = 1.3 \times 10^{-3} n_7^{-1}$$



- Condition for an outflow:



- Condition for an outflow:  $T_0 > T_{\text{vir},r}(r_0) \Gamma_c^{1/4}$

- Condition for an outflow:  $T_0 > T_{\text{vir,r}}(r_0) \Gamma_c^{1/4}$

- Virial temperature in radiatively-dominated plasma:



- Condition for an outflow:  $T_0 > T_{\text{vir,r}}(r_0) \Gamma_c^{1/4}$

- Virial temperature in radiatively-dominated plasma:

$$T_{\text{vir,r}} = \left( \frac{GM\rho}{ar} \right)^{1/4} \simeq 312 \left( \frac{n_5 M_7}{r_{\text{pc}}} \right)^{1/4} - 987 \left( \frac{n_7 M_7}{r_{\text{pc}}} \right)^{1/4} \text{ K}$$

$$g_{\text{rad}} = \kappa F / c = \kappa \lambda dE / dl \sim (1 / \rho) dE / dl \sim \frac{GM}{r^2}$$



$$\lambda = 1/(\kappa\rho)$$

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

$$\lambda = 1/(\kappa\rho)$$

- mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$



$\lambda = 1/(\kappa\rho)$  - mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:

$\lambda = 1/(\kappa\rho)$  - mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:  $dE/dl \sim E/L \sim E/R$



$\lambda = 1/(\kappa\rho)$  - mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:  $dE/dl \sim E/L \sim E/R$

$$T_{\text{vir,r}} = \left(\frac{GM\rho}{ar}\right)^{1/4} \simeq 312 \left(\frac{n_5 M_7}{r_{\text{pc}}}\right)^{1/4} - 987 \left(\frac{n_7 M_7}{r_{\text{pc}}}\right)^{1/4} \text{ K}$$

$$\lambda = 1/(\kappa\rho)$$

- mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:  $dE/dl \sim E/L \sim E/R$

$$T_{\text{vir,r}} = \left( \frac{GM\rho}{ar} \right)^{1/4} \simeq 312 \left( \frac{n_5 M_7}{r_{\text{pc}}} \right)^{1/4} - 987 \left( \frac{n_7 M_7}{r_{\text{pc}}} \right)^{1/4} \text{ K}$$



$\lambda = 1/(\kappa\rho)$  - mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:  $dE/dl \sim E/L \sim E/R$

$$T_{\text{vir,r}} = \left( \frac{GM\rho}{ar} \right)^{1/4} \simeq 312 \left( \frac{n_5 M_7}{r_{\text{pc}}} \right)^{1/4} - 987 \left( \frac{n_7 M_7}{r_{\text{pc}}} \right)^{1/4} \text{ K}$$

Optically thin transport:

$\lambda = 1/(\kappa\rho)$  - mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:  $dE/dl \sim E/L \sim E/R$

$$T_{\text{vir,r}} = \left( \frac{GM\rho}{ar} \right)^{1/4} \simeq 312 \left( \frac{n_5 M_7}{r_{\text{pc}}} \right)^{1/4} - 987 \left( \frac{n_7 M_7}{r_{\text{pc}}} \right)^{1/4} \text{ K}$$

Optically thin transport:  $dE/dl \sim E/\lambda$



$\lambda = 1/(\kappa\rho)$  - mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:  $dE/dl \sim E/L \sim E/R$

$$T_{\text{vir,r}} = \left( \frac{GM\rho}{ar} \right)^{1/4} \simeq 312 \left( \frac{n_5 M_7}{r_{\text{pc}}} \right)^{1/4} - 987 \left( \frac{n_7 M_7}{r_{\text{pc}}} \right)^{1/4} \text{ K}$$

Optically thin transport:  $dE/dl \sim E/\lambda$

$$T_{\text{vir,flx}} = \left( \frac{4GM}{ar^2\kappa} \right)^{1/4} = 292 \left( \frac{M_7}{r_{\text{pc}}^2 \kappa_{10}} \right)^{1/4} \text{ K}$$

$\lambda = 1/(\kappa\rho)$  - mean free path

$$g_{\text{rad}} = \kappa F/c = \kappa\lambda dE/dl \sim (1/\rho) dE/dl \sim \frac{GM}{r^2}$$

Optically thick diffusion:  $dE/dl \sim E/L \sim E/R$

$$T_{\text{vir,r}} = \left( \frac{GM\rho}{ar} \right)^{1/4} \simeq 312 \left( \frac{n_5 M_7}{r_{\text{pc}}} \right)^{1/4} - 987 \left( \frac{n_7 M_7}{r_{\text{pc}}} \right)^{1/4} \text{ K}$$

Optically thin transport:  $dE/dl \sim E/\lambda$

$$T_{\text{vir,fx}} = \left( \frac{4GM}{ar^2\kappa} \right)^{1/4} = 292 \left( \frac{M_7}{r_{\text{pc}}^2 \kappa_{10}} \right)^{1/4} \text{ K}$$



- $\dot{M}$  as an eigenvalue problem
- Mass of the “torus”
- Speed of an outflow
- AGN feedback

# RADIATION HYDRODYNAMICS OF AGN OBSCURATION



# RADIATION HYDRODYNAMICS OF AGN OBSCURATION


$$D_t \rho + \rho \nabla \cdot \mathbf{v} = 0,$$

$$D_t \mathbf{v} = -\frac{1}{\rho} \nabla \mathbf{p} + \mathbf{g}_{\text{rad}} - \nabla \Phi,$$

# RADIATION HYDRODYNAMICS OF AGN OBSCURATION

$$D_t \rho + \rho \nabla \cdot \mathbf{v} = 0,$$

$$D_t \mathbf{v} = -\frac{1}{\rho} \nabla \mathbf{p} + \mathbf{g}_{\text{rad}} - \nabla \Phi,$$

$$\frac{\mathbf{F}}{c} \kappa_d$$




# RADIATION HYDRODYNAMICS OF AGN OBSCURATION

$$D_t \rho + \rho \nabla \cdot \mathbf{v} = 0,$$

$$\frac{\mathbf{F}}{c} \kappa_d$$

$$D_t \mathbf{v} = -\frac{1}{\rho} \nabla \mathbf{p} + \mathbf{g}_{\text{rad}} - \nabla \Phi,$$

$$\rho D_t \left( \frac{e}{\rho} \right) = -p \nabla \cdot \mathbf{v} - 4\pi \chi_P B + c \chi_E E,$$

$$\rho D_t \left( \frac{E}{\rho} \right) = -\nabla \cdot \mathbf{F} - \nabla \mathbf{v} : \mathbf{P} + 4\pi \chi_P B - c \chi_E E$$

# RADIATION HYDRODYNAMICS OF AGN OBSCURATION

$$D_t \rho + \rho \nabla \cdot \mathbf{v} = 0,$$

$$D_t \mathbf{v} = -\frac{1}{\rho} \nabla \mathbf{p} + \mathbf{g}_{\text{rad}} - \nabla \Phi,$$

$$\rho D_t \left( \frac{e}{\rho} \right) = -p \nabla \cdot \mathbf{v} - 4\pi \chi_P B + c \chi_E E,$$

$$\rho D_t \left( \frac{E}{\rho} \right) = -\nabla \cdot \mathbf{F} - \nabla \mathbf{v} : \mathbf{P} + 4\pi \chi_P B - c \chi_E E$$

$$\frac{\mathbf{F}}{c} \kappa_d$$

Radiation Energy Density

$$E(\mathbf{r}, \mathbf{t}) = \frac{1}{c} \int_0^\infty \int d\Omega \mathbf{I}(\mathbf{r}, \Omega, \nu, \mathbf{t})$$



# RADIATION HYDRODYNAMICS OF AGN OBSCURATION

$$D_t \rho + \rho \nabla \cdot \mathbf{v} = 0,$$

$$D_t \mathbf{v} = -\frac{1}{\rho} \nabla p + \mathbf{g}_{\text{rad}} - \nabla \Phi,$$

$$\rho D_t \left( \frac{e}{\rho} \right) = -p \nabla \cdot \mathbf{v} - 4\pi \chi_P B + c \chi_E E,$$

$$\rho D_t \left( \frac{E}{\rho} \right) = -\nabla \cdot \mathbf{F} - \nabla \mathbf{v} : \mathbf{P} + 4\pi \chi_P B - c \chi_E E$$

$$\frac{\mathbf{F}}{c} \kappa_d$$

Radiation Energy Density

$$E(\mathbf{r}, \mathbf{t}) = \frac{1}{c} \int_0^\infty \int d\Omega \mathbf{I}(\mathbf{r}, \Omega, \nu, \mathbf{t})$$

Radiation Flux

$$\mathbf{F}(\mathbf{r}, \mathbf{t}) = \int_0^\infty \int d\Omega \hat{\mathbf{n}} \mathbf{I}(\mathbf{r}, \Omega, \nu, \mathbf{t})$$

# RADIATION HYDRODYNAMICS OF AGN OBSCURATION

$$D_t \rho + \rho \nabla \cdot \mathbf{v} = 0,$$

$$D_t \mathbf{v} = -\frac{1}{\rho} \nabla p + \mathbf{g}_{\text{rad}} - \nabla \Phi,$$

$$\rho D_t \left( \frac{e}{\rho} \right) = -p \nabla \cdot \mathbf{v} - 4\pi \chi_P B + c \chi_E E,$$

$$\rho D_t \left( \frac{E}{\rho} \right) = -\nabla \cdot \mathbf{F} - \nabla \mathbf{v} : \mathbf{P} + 4\pi \chi_P B - c \chi_E E$$

$$\frac{\mathbf{F}}{c} \kappa_d$$

Radiation Energy Density

$$E(\mathbf{r}, t) = \frac{1}{c} \int_0^\infty \oint d\Omega \mathbf{I}(\mathbf{r}, \Omega, \nu, t)$$

Radiation Flux

$$\mathbf{F}(\mathbf{r}, t) = \int_0^\infty \oint d\Omega \hat{\mathbf{n}} \mathbf{I}(\mathbf{r}, \Omega, \nu, t)$$

Radiation Pressure Tensor

$$\mathbf{P}(\mathbf{r}, t) = \frac{1}{c} \int_0^\infty \oint d\Omega \hat{\mathbf{n}} \hat{\mathbf{n}} I(\mathbf{r}, \Omega, \nu, t)$$





---

Dorodnitsyn et al. , ApJ, 2011, 741, 29  
Dorodnitsyn et al., ApJ, 2012, 747, 8



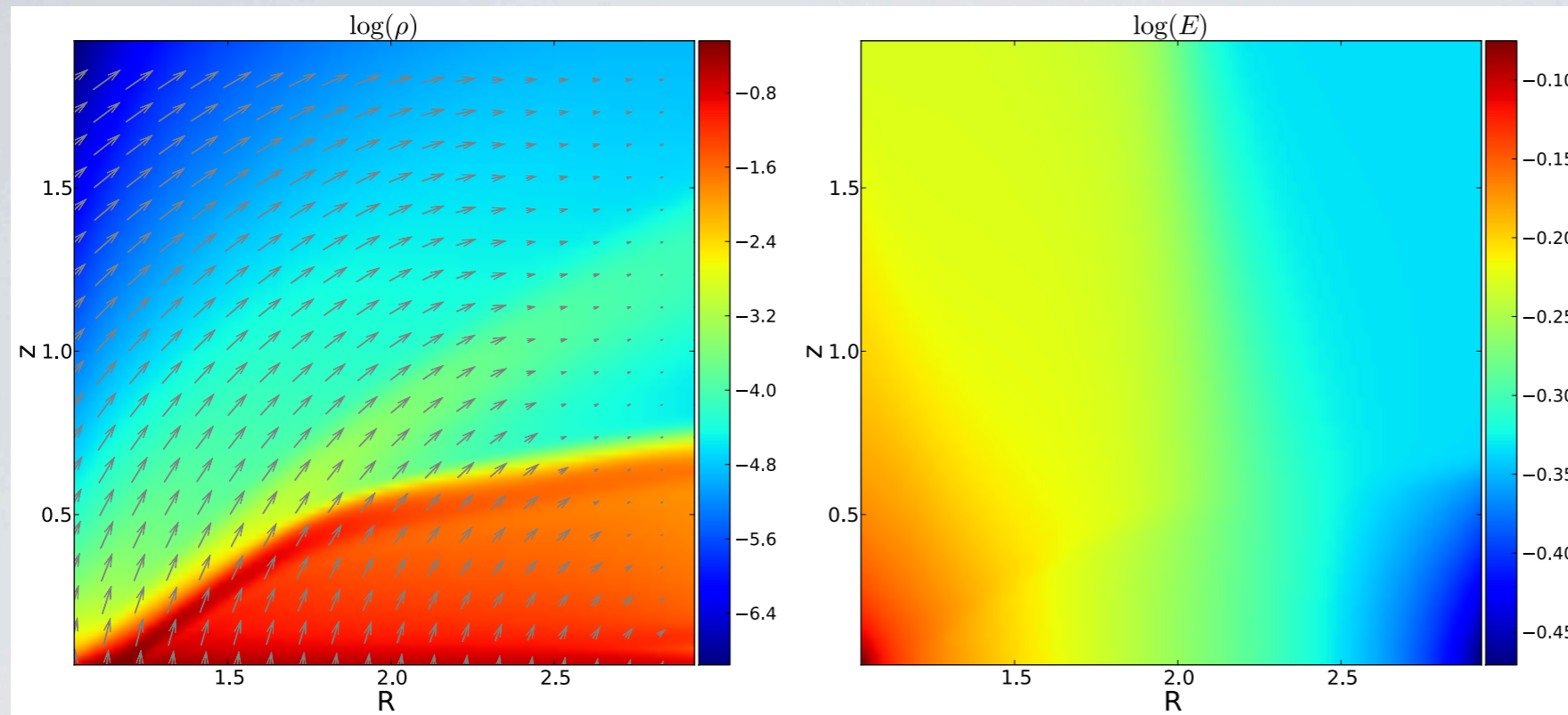
---

Dorodnitsyn et al. , ApJ, 2011, 741, 29  
Dorodnitsyn et al., ApJ, 2012, 747, 8

$$L = 0.5 L_{\text{edd}}$$

$$R = 1 \text{ pc}$$

$$\tau_{\text{T}} = 0.53$$




---

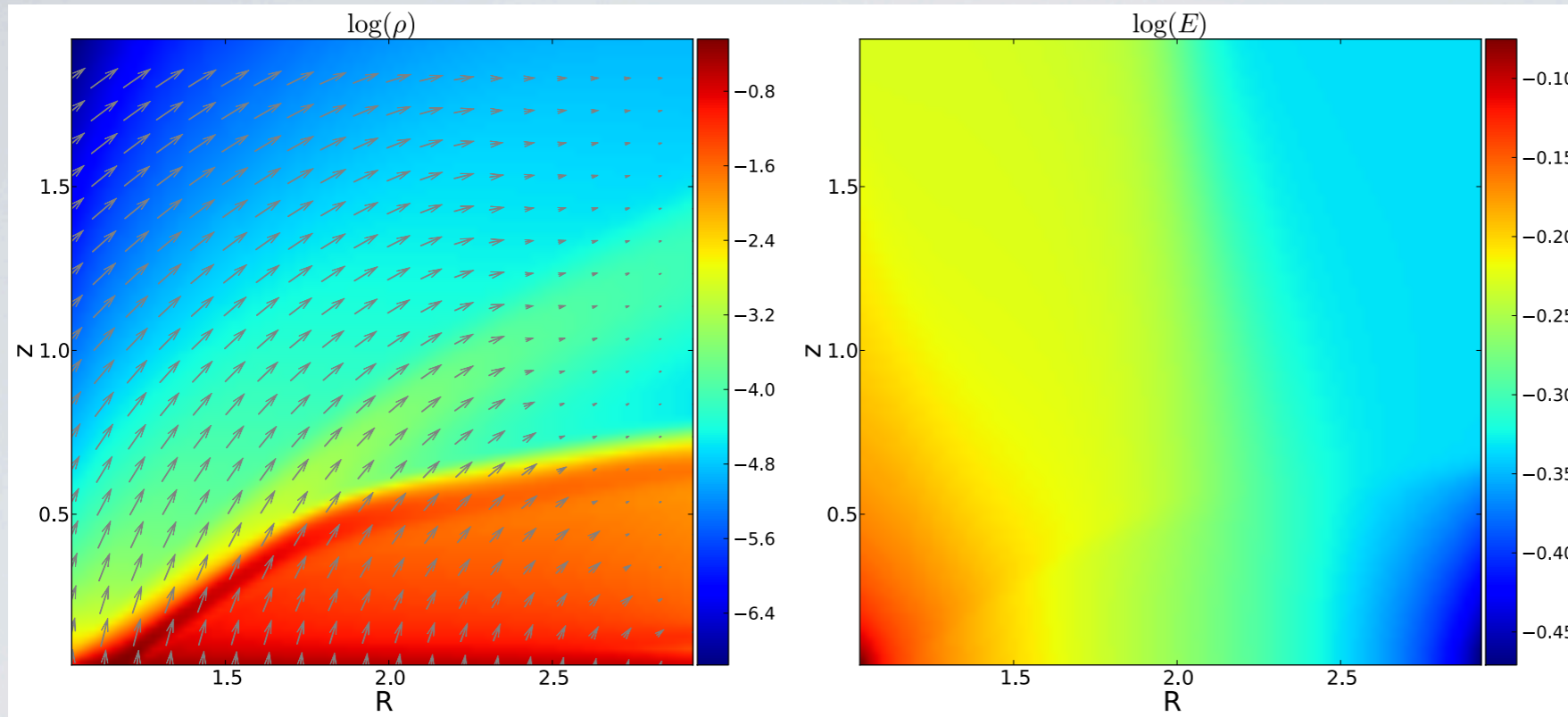
Dorodnitsyn et al. , ApJ, 2011, 741, 29  
Dorodnitsyn et al., ApJ, 2012, 747, 8

$$L = 0.5 L_{\text{ed}}d$$

$$R = 1 \text{ pc}$$

$$\tau_{\text{T}} = 0.53$$



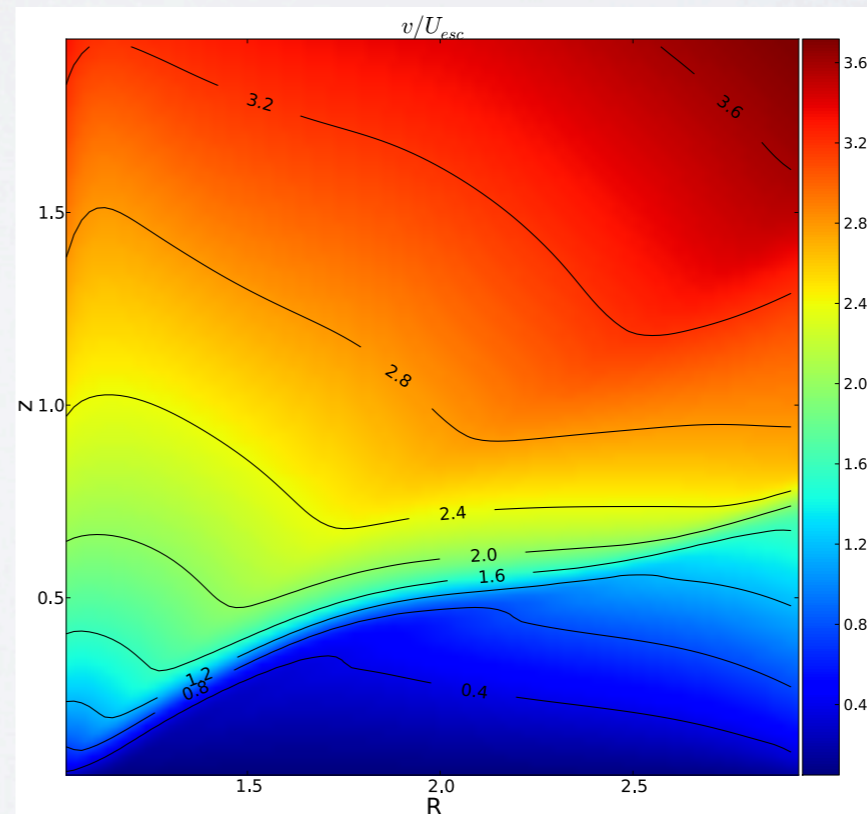
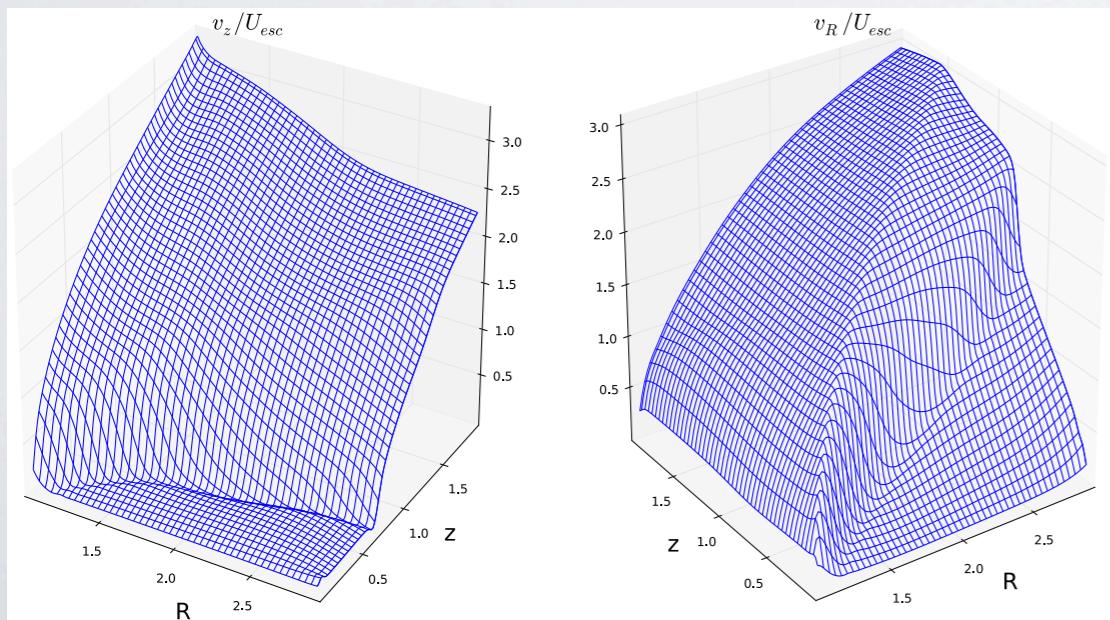


Dorodnitsyn et al. , ApJ, 2011, 741, 29  
 Dorodnitsyn et al., ApJ, 2012, 747, 8

$$L = 0.5 L_{\text{edd}}$$

$$R = 1 \text{ pc}$$

$$\tau_T = 0.53$$



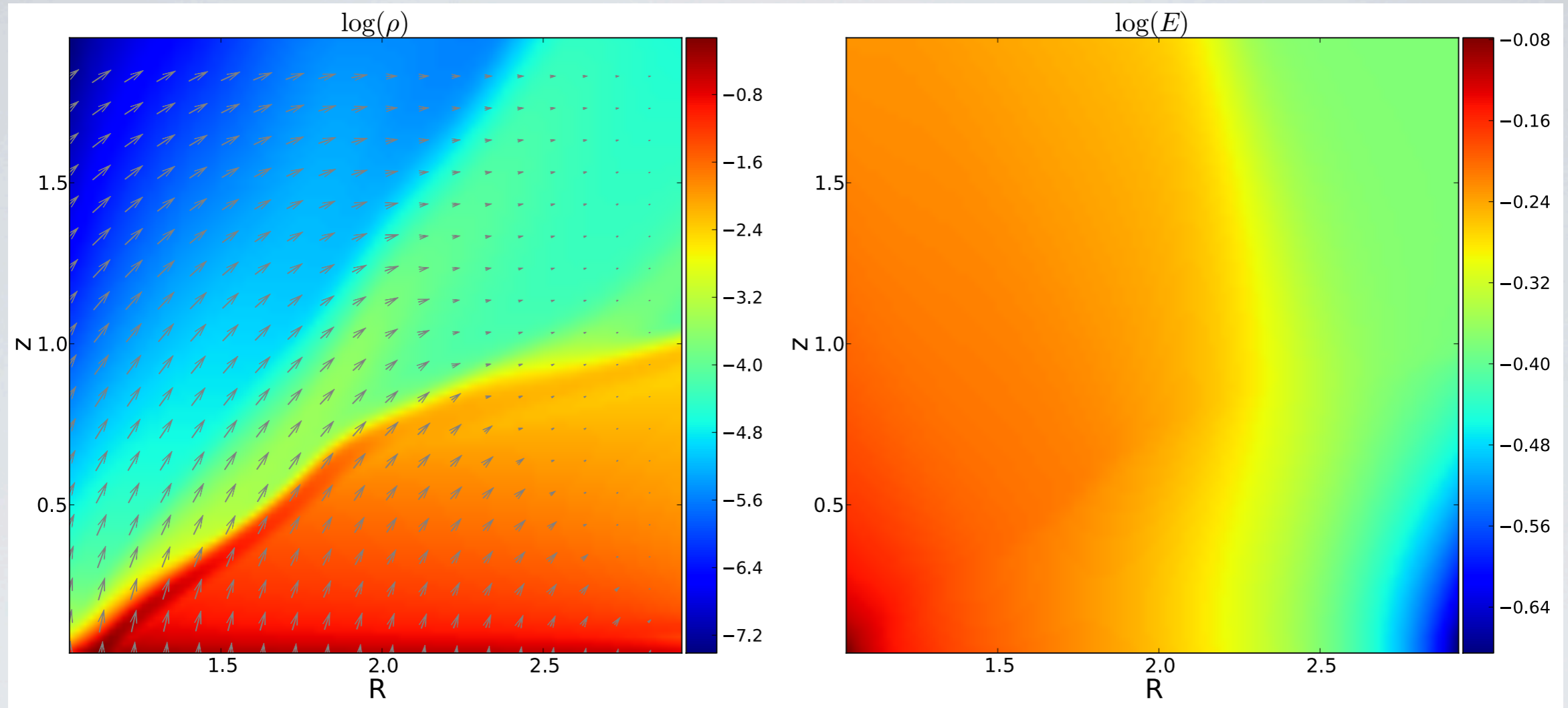




$$L = 0.05 L_{\text{edd}}$$

$$R = 1 \text{ pc}$$

$$\tau_{\text{T}} = 1.8$$

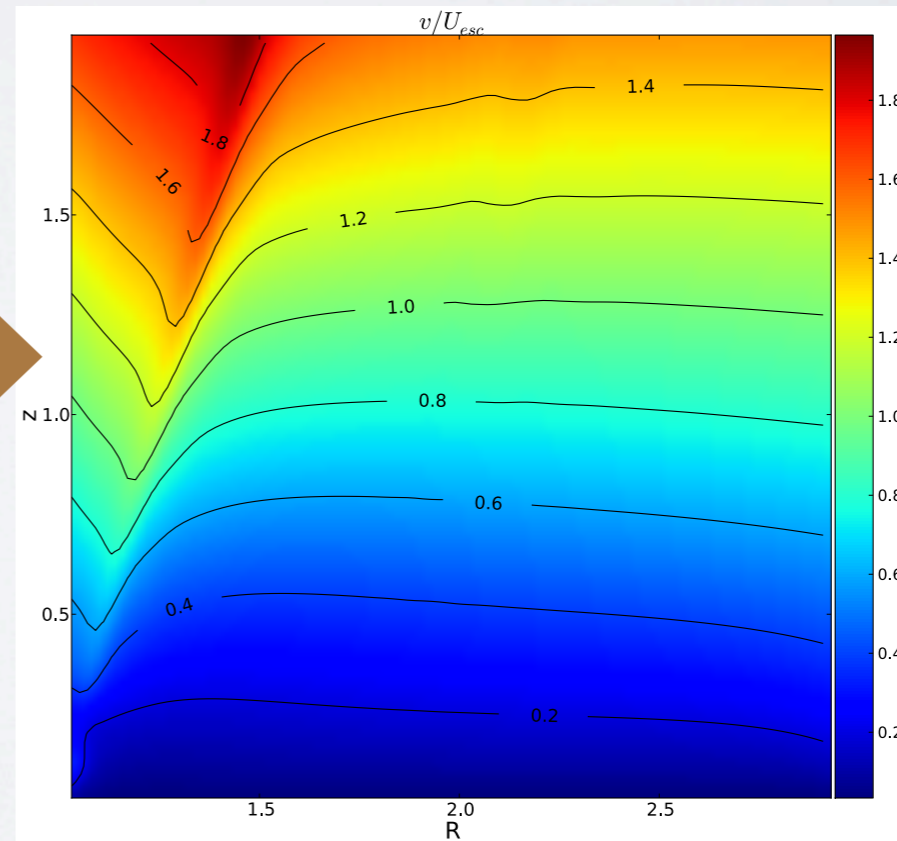
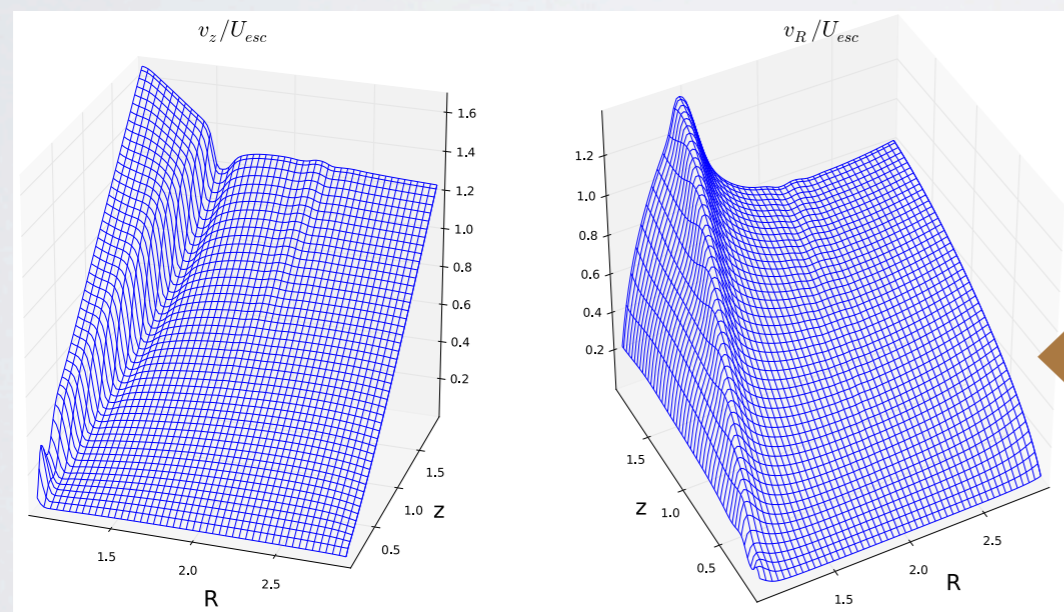
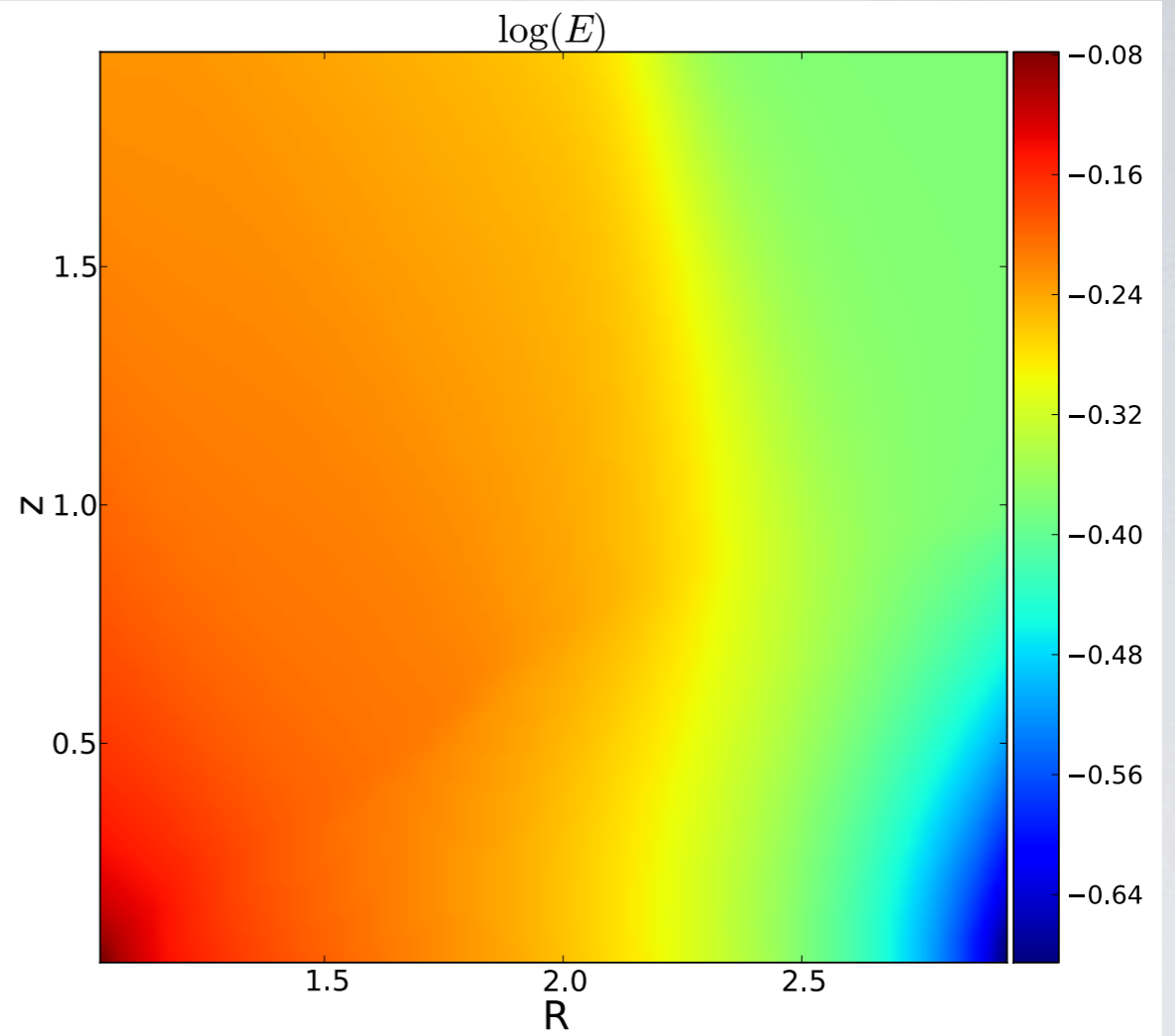
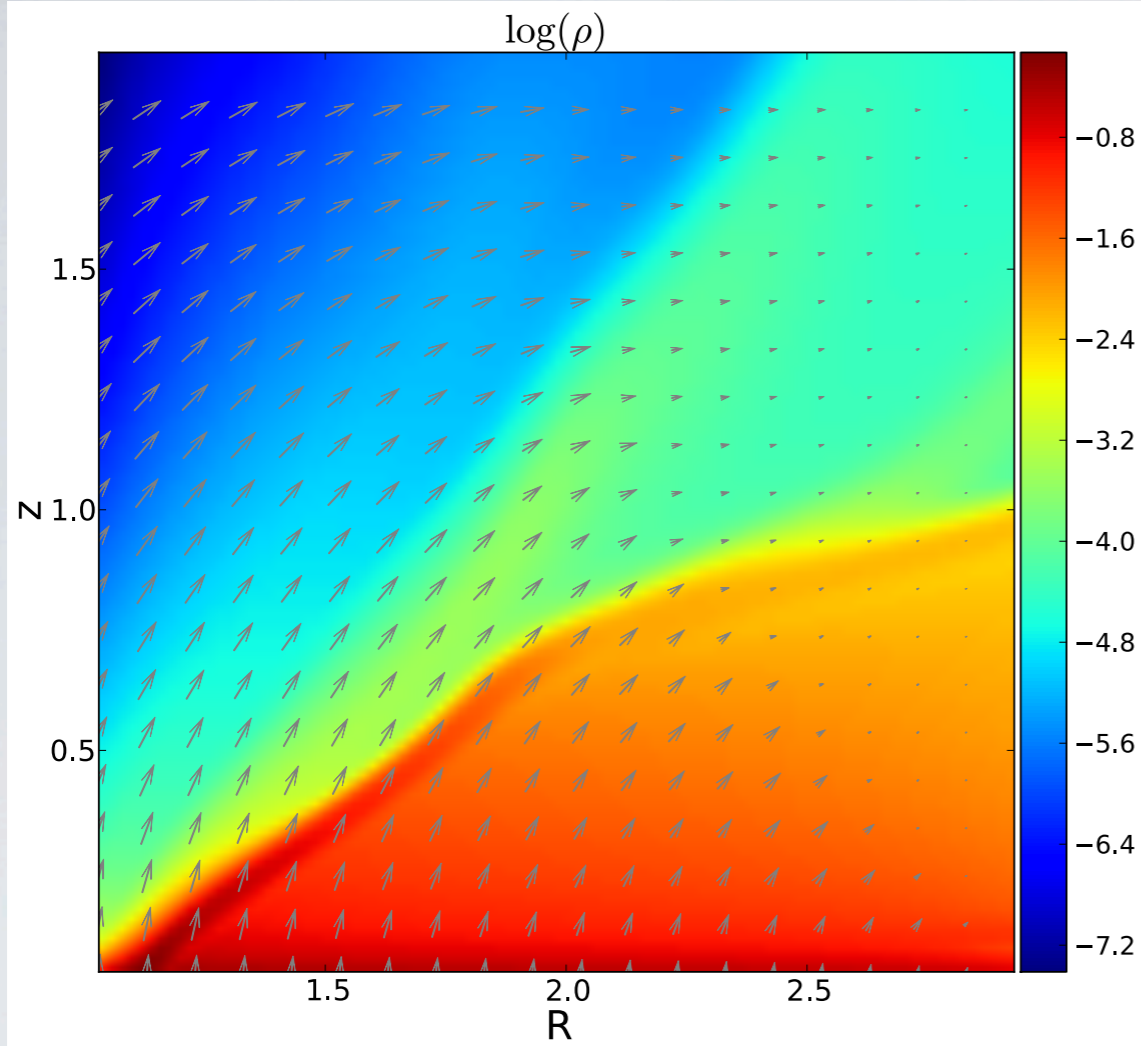


$$L = 0.05 L_{\text{edd}}$$

$$R = 1 \text{ pc}$$

$$\tau_{\text{T}} = 1.8$$





$$L = 0.05 L_{\text{edd}}$$

$$R = 1 \text{ pc}$$

$$\tau_{\text{T}} = 1.8$$

- As the luminosity of the central machine exceeds  $0.1 L_{\text{edg}}$
- The geometrically thick wind is formed
- This “wind” is obscuring AGN at high inclinations

RHD, time-dependent 2.5D

Model	$\Gamma$	$R_0$	$\tau_T$	$n_0$	$\langle v \rangle$	$v_{\text{max}}^*$	$L_{\text{kin}}$	$L_{\text{bol}}$	$\dot{M}$
1	0.5	1	0.53	$3 \cdot 10^5$	138.4	666.61	$1.23 \cdot 10^{41}$	$6.24 \cdot 10^{44}$	1.71
2	0.3	1	0.53	$3 \cdot 10^5$	104.87	560.27	$6.38 \cdot 10^{40}$	$3.74 \cdot 10^{44}$	1.38
3	0.1	1	0.53	$3 \cdot 10^5$	55.82	373.46	$1.48 \cdot 10^{40}$	$1.24 \cdot 10^{44}$	1.85
4	0.05	1	0.53	$3 \cdot 10^5$	36.63	282.23	$6 \cdot 10^{39}$	$6.25 \cdot 10^{43}$	0.64
5	0.8	1	1.8	$1 \cdot 10^6$	112.62	765	$4.11 \cdot 10^{41}$	$9.99 \cdot 10^{44}$	5
6	0.5	1	1.8	$1 \cdot 10^6$	90.55	513.42	$2.56 \cdot 10^{41}$	$6.24 \cdot 10^{44}$	4.18
7	0.3	1	1.8	$1 \cdot 10^6$	73.18	348.45	$1.37 \cdot 10^{41}$	$3.74 \cdot 10^{44}$	3.37
8	0.1	1	1.8	$1 \cdot 10^6$	59.73	159.74	$3.09 \cdot 10^{40}$	$1.24 \cdot 10^{44}$	1.88
9	0.05	1	1.8	$1 \cdot 10^6$	42.8	129.39	$1.31 \cdot 10^{40}$	$6.25 \cdot 10^{43}$	1.39



# CONCLUSIONS

# STATIC TORUS

# CONCLUSIONS



# CONCLUSIONS

“Dynamic Torus”

# CONCLUSIONS

- Active Galactic Nuclei (AGN), Seyfert galaxies and quasars, are powered by luminous accretion and often accompanied by winds which are powerful enough to affect the AGN mass budget, and whose observational appearance bears an imprint of processes which are happening within the central parsec around the black hole (BH).
- Our results demonstrate that for AGN luminosities greater than  $0.1 L_{\text{edd}}$  external illumination can support a geometrically thick obscuration via outflows driven by the infrared radiation pressure.
- The terminal velocity of marginally Compton-thin models  $0.2 < \tau_{\text{T}} < 0.6$  is comparable or greater than the escape velocity.
- In Compton thick models the maximum value of the vertical component of the velocity is lower than the escape velocity suggesting that a significant part of our torus is in the form of failed

Pressure of the infrared radiation on dust is crucial. Dusty winds may or may not be failed, i.e. returning. This flow is known aka “AGN torus”