# Asymptotic Accuracy of the Equilibrium-Diffusion Approximation

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Asymptotic Accuracy

## Overview

#### Introduction

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- Non-relativistic equilibrium-diffusion is a fundamental limit of non-relativistic radiation-hydrodynamics.
- It is associated with optically-thick media, strong radiation-material coupling, and non-relativistic material speeds.
- An asymptotic expansion of the non-relativistic radiation-hydrodynamic (NRRH) solution has previously been shown to satisfy the non-relativistic equilibrium-diffusion (NRED) equations to leading-order (Lowrie, Morel, and Hittinger).
- In the current work, we have shown that this asymptotic expansion of the NRRH solution satisfies the NRED equations to first-order with an  $O(\epsilon^2)$  error.

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#### The Asymptotic Procedure

- Start with NRRH equations coupled to fully relativistic lab-frame transport equation and non-dimensionalize these equations.
- Identify non-dimensional parameters of interest and scale these by positive or negative powers of a small parameter ε:

$$\mathcal{U} \equiv rac{\mathbf{V}_{\infty}}{\mathbf{c}} \sim \mathcal{O}(\varepsilon)$$
  
 $\mathcal{R} \equiv rac{
ho_{\infty} u_{\infty}^2}{\mathbf{a}_r T_{\infty}^4} \sim \mathcal{O}(1)$   
 $\mathcal{L} \equiv rac{\ell_{\infty}}{\lambda_{t,\infty}} \sim \mathcal{O}(\varepsilon^{-1})$   
 $\mathcal{L}_{\mathbf{s}} \equiv rac{\lambda_{t,\infty}}{\lambda_{\mathbf{s},\infty}} \sim \mathcal{O}(\varepsilon).$ 

## The Asymptotic Procedure

- Equate terms multiplying each power of  $\varepsilon$  and manipulate the resulting equations to find closed equations for the unknowns in the power-series expansion.
- The first-order expansion to the NRRH solution, U<sub>0</sub> + εU<sub>1</sub>, satisfies the NRED equations, but the second-order expansion does not.

## Why This is Important

- Radiation-hydrodynamics models and discretizations should preserve the asymptotic accuracy of the NRRH equations in the NRED limit.
- Discretizations and models that are not asymptotic preserving to leading-order are usually impractical for calculations that are strongly asymptotic due to the need for arbitrarily excessive mesh resolution.
- Discretizations and models that are partially asymptotic preserving will not be as accurate as they should be for calculations that are strongly asymptotic.

## Why This is Important

• For instance, the Wilson flux limiter for radiative diffusion,

$$D = 1/(3\sigma_t + \|\overrightarrow{
abla} \mathcal{E}\|/\mathcal{E})$$

is known to limit too strongly.

- This can be explained by the fact that it preserves the NRED limit to leading order but not to first order, i.e., the solution departs from diffusion with order ε instead of order ε<sup>2</sup> (Morel).
- In contrast, the Larsen limiter,

$$D = 1/\sqrt{(9\sigma_t^2 + [\| \overrightarrow{
abla} \mathcal{E} \| / \mathcal{E}]^2)}$$

preserves the NRED limit to first-order and is more accurate in general than the Wilson limiter (Morel).

## **Comoving-Frame Diffusion Model**

- Start with comoving-frame energy and momentum equations.
- Assume  $\mathcal{P}_{r,0} = \frac{1}{3}\mathcal{E}_{r,0}$ .
- Obtain Fick's Law by setting certain terms to zero.

$$\partial_t \mathcal{E}_0 - \partial_i \left( \frac{c}{3\sigma_t} \partial_i \mathcal{E}_0 \right) + \partial_i \left( \mathbf{v}_i \mathcal{E}_0 \right) + \frac{1}{3} \mathcal{E}_0 \partial_i \mathbf{v}_i = c \sigma_a \left( a T^4 - \mathcal{E}_0 \right) \quad ,$$
$$\mathcal{F}_{0,i} = -\frac{c}{3\sigma_t} \partial_i \mathcal{E}_0 \quad .$$

### Lab-Frame Diffusion Model

- Start with lab-frame energy and momentum equations.
- Assume  $\mathcal{P}_r = \frac{1}{3}\mathcal{E}_r$ .
- Set the time-derivative of the radiation flux to zero.

$$\partial_{t} \mathcal{E} - \partial_{i} \frac{c}{3\sigma_{t}} \partial_{i} \mathcal{E} + \partial_{i} \left\{ v_{i} \left[ \frac{4}{3} \mathcal{E} + \frac{\sigma_{a}}{\sigma_{t}} \left( aT^{4} - \mathcal{E}_{0} \right) \right] \right\} = c\sigma_{a} \left( aT^{4} - \mathcal{E}_{0} \right) - \frac{\sigma_{t}}{c} v_{i} \mathcal{F}_{0,i} .$$

$$\mathcal{F}_{i} = -\frac{c}{3\sigma_{t}} \partial_{i} \mathcal{E} + v_{i} \frac{4}{3} \mathcal{E} + v_{i} \frac{\sigma_{a}}{\sigma_{t}} \left( aT^{4} - \mathcal{E}_{0} \right) .$$

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## Comparison of Model Accuracy

- The energy and momentum equations in the comoving- and lab-frames are equivalent, i.e., they transform into one another neglecting higher-order v<sub>i</sub>/c terms.
- The diffusion approximations in the comoving- and lab-frames are fundamentally different, i.e., they do not transform into one another.
- It has been suggested that the lab-frame approximation should be fundamentally flawed because the radiation intensity *does not* become isotropic in the lab-frame.
- However, our asymptotic analyses show that both approximations, when coupled to the Euler equations, preserve the NRED limit to first order.
- This indicates that they are equally valid in the NRED limit.

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#### Conclusions

- Knowledge of the asymptotic accuracy of the NRED approximation enable us to fully evaluate the accuracy of radiation-hydrodynamics models and discretizations in the NRED limit.
- Failure to preserve the accuracy of the NRED limit can lead to models and discretizations with significant flaws.